

Problem Set 4: due Tuesday, February 27, 2018

Useful references:

Frank Shu; "Gas Dynamics"

P. Drazin and W. Reid; "Hydrodynamic Stability"

And, of course, Landau and Lifshitz

1) Tapas

Each of these questions does not require more than a few lines of calculation. *Don't* make them longer than they need be. Do state your reasoning clearly. Try to do these closed book.

- a) What is the width of the laminar boundary layer at the bottom of a viscous fluid rotating at Ω ?
- b) What is the width of a laminar boundary layer of a stagnation flow?
- c) What shaped eddy is most effective at Rayleigh–Bernard convection when $\Omega^2 \gg g\alpha\beta$?
- d) How will the strength of a vortex tube evolve when the fluid density within rises linearly in time?
- e) A sphere of radius R moves thru an inviscid fluid line (i.e. obeys Euler, not Navier–Stokes) with a free surface. The sphere moves at V , at various depths d . What happens? Estimate the drag on the sphere.

2) Calculate the flow in a boundary layer in a converging channel between two non-parallel plates. See Landau sections 23, 39. Explain your similarity solution carefully.

3) Now consider a rotating fluid which is also compressible and *self-gravitating*. For the latter, include a body force $\underline{f} = -\nabla\phi$ where:

$$\nabla^2\phi = 4\pi G\rho$$

Take $\underline{\Omega} = \Omega\hat{z}$, as usual.

a) For $\underline{k} = k\hat{z}$, show

$$\omega^2 = k^2 c_s^2 - 4\pi G\rho_0.$$

Welcome to the Jeans instability! What might be the significance of the marginally stable length scale?

- b) For $k = k\hat{x}$, show:

$$\omega^2 = k^2 c_s^2 - 4\pi G \rho_0 + 4\Omega^2.$$

When are all modes stable?

- c) Why might this result be of interest in the context of galaxy structure?
 d) How might Jeans instabilities evolve nonlinearly? Why might cooling effects be important here?

- 4) Consider a sheared flow $\underline{v}(z)\hat{x}$ in a stably stratified fluid with $\underline{g} = -g\hat{z}$.

- a) Derive the 2D wave eigenmode equation, called the Taylor–Goldstein equation. This extends the Rayleigh equation from the last problem of Set 3.
 b) As before, take $\omega = \omega_r + i\gamma$, and substitute $H = \phi/(v-c)^{1/2}$ where c is the along-stream phase velocity. Multiply the re-scaled Taylor–Goldstein equation by H^* and derive a quadratic form.
 c) Show for $\gamma > 0$, the *Richardson number* must satisfy:

$JN^2/V'^2 < 1/4$, for shear-driven instability.

Here $\frac{JN^2}{V'^2} = \frac{-g}{\rho} \frac{d\rho}{dz} / V'^2$, is the Richardson number.

- d) What is the physics of the Richardson number criterion? What does it mean, in simple terms?
 e) Extra Credit: Can you give a grunt-and-hand gesture derivation of the Richardson number criterion, based on simple physical ideas of competition of energetics?
- 5) Calculate the boundary layer thickness and flow field near a semi-infinite plate, assuming laminar flow. This is the Blasius problem. Follow the self-similarity approach discussed in Landau and Lifshitz. Get as far as you can with the numerical factors.

- 6a) Using Oseen's approximation, calculate the shape, flow field and pressure field for the laminar wake, assuming conical symmetry, as discussed in class.
- b) Now, do the same for a turbulent wake, exploiting the idea of a turbulent viscosity, as introduced in class.
- c) How far behind the object does the wake transition from turbulent to laminar?